GINA – A Polarized Neutron Reflectometer at the Budapest Neutron Centre

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The setup, capabilities and operation parameters of the neutron reflectometer GINA, the recently installed "Grazing Incidence Neutron Apparatus" at the Budapest Neutron Centre, are introduced. GINA, a dance-floor-type, constant-energy, angle-dispersive reflectometer is equipped with a 2D position-sensitive detector to study specular and off-specular scattering. Wavelength options between 3.2 and 5.7 Å are available for unpolarized and polarized neutrons. Spin polarization and analysis are achieved by magnetized transmission supermirrors and radio-frequency adiabatic spin flippers. As a result of vertical focusing by a five-element pyrolytic graphite monochromator, the reflected intensity from a $20 \times 20 \text{ mm}^2$ sample has been doubled. GINA is dedicated to studies of magnetic films and heterostructures, but unpolarized options for non-magnetic films, membranes, and other surfaces are also provided. Shortly after its startup, reflectivity values as low as 9×10^{-6} have been measured by the instrument. The instrument capabilities are demonstrated by a non-polarized and a polarized reflectivity experiment on a Si wafer and on a magnetic film of $\begin{bmatrix} 62 \text{Ni}/\text{nat}\text{Ni} \end{bmatrix}_5$ isotope-periodic layer composition. The facility is now open for the international user community. Its further development is underway establishing new sample environment options and spin analysis of off-specularly scattered radiation as well as further decreasing the background.

I. INTRODUCTION

A. Neutron reflectometry in materials nanoscience

The ever increasing need for product advancement and miniaturization keeps membranes, thin film assemblies, magnetic and non-magnetic multilayer and patterned heterostructures in the limelight of materials science and technological development. In recent years more and more complex methods, instrumental and evaluation types have emerged to meet the new challenges. Due to the matching wavelength range of cold neutrons and their extreme sensitivity to the interface structure and to the internal magnetic fields, neutron reflectometry (NR) is a rather capable non-destructive method to investigate such nanostructures. Reflected intensity measured as a function of the momentum transfer Q_z , perpendicular to the sample surface provides information on the scattering length density (SLD) depth profile. Normalized reflectivities are usually recorded from unity in the total reflection regime to Q_z values of about 0.2 Å⁻¹, where it drops by five to six orders of magnitude. Layer thicknesses appear as regular features in the reflected intensity and may be modeled using optical formalisms.^{1,2} Similar but spin dependent formalisms apply for polarized neutron reflectometry (PNR) in magnetic studies.^{3–5} By fitting the appropriate model to the measured intensities as a function of Q_z , one can extract layer thickness, interfacial roughness, and depth-dependent SLD, as well as magnetization

profile. By measuring scattered intensity as a function of an in-plane wave vector component Q_x , one can, in addition, characterize lateral structures.⁶ In the lateral direction, the interface may be rough on different length scales or may display a defined periodicity, resulting in diffuse scattering or Bragg reflections in Q_x scans, respectively. Polarized neutron reflectometry (PNR) not only provides an isotope-selective atomic density depth profile (as well in the case of deeply buried layers) with a spatial resolution of a few nanometers,^{7,8} but it is also a highly sensitive magnetometry to determine the vector properties of the magnetization. The prototype polarized neutron reflectometer was developed and built at Argonne in the 1980s.^{9,10} The increased interest in magnetic thin film analytical instruments triggered by the discovery of the giant magnetoresistance and related phenomena^{11,12} resulted in a boom of PNR studies^{13–15} as well as the construction of a number of new neutron reflectometers with polarization option at neutron sources across the globe. Here we report on the design, the construction and the operation parameters of the "Grazing Incidence Neutron Apparatus" (GINA)¹⁶ a recently installed neutron reflectometer at the Budapest Research Reactor $(BRR)^{17}$ of the Budapest Neutron Centre (BNC), Hungary.

B. The neutron source at the Budapest Neutron Centre

The research reactor at BNC is of tank type, it is moderated and cooled by light water. The area of the reactor hall, the cold neutron source (CNS) guide hall (for the cold neutron instruments) and the second guide hall (for a thermal time of flight (TOF) instrument) are approxi-

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FIG. 1. The layout of the CNS guide hall with the GINA instrument (IBMS: In-beam Mössbauer spectrometer; SANS: Small angle neutron scattering; REF: Neutron reflectometer; PGAA:.Promt gamma activation analysis; NIPS: Neutron induced prompt gamma spectrometer; ATHOS: Tripleaxis spectrometer)

mately 600, 400 and 120 m^2 , respectively. Until the year 2009 the reactor was fuelled by 36% enriched 235 U. The core, which is surrounded by a solid beryllium reflector, is being converted and will be fully converted to 20% enriched Russian uranium fuel of type VVR-M2 by the end of 2012. The thermal power, the mean power density, the maximum thermal neutron flux and the maximum cooling water outlet temperature are 10 MW, 39.7 W/m^3 , $2.1 \times 10^{14} \mathrm{n/cm^{2}s}$ and $60^{\circ}\mathrm{C}$, respectively. One reactor cycle at BNC lasts 10 effective days, which is followed by a short break for a weekend. The number of reactoroperation days per year varied from 156 to 165 in the recent years and similar figures are foreseen for the coming years.¹⁷ The reactor has eight radial and two tangential beam tubes. On the tangential beam tube 10 the cold neutron source is installed. Three neutron guides (10/1,10/2 and 10/3) originating at the CNS provide cold neutrons to the facilities in the CNS guide hall (see Fig. 1). REF, the first neutron reflectometer at BNC has recently been moved from cold guide 10/3 to cold guide 10/1 and is primarily used for quality test of neutron mirrors and other optical elements produced by Mirrotron Ltd., Budapest. The new GINA reflectometer is installed in the former position of REF on the cold guide 10/3, closest to the wall of the reactor building. The GINA facility occupies approx. 60 square meters including its well polished marble "dance floor" (25 m^2) and its control hutch.

II. GENERAL OVERVIEW OF THE GINA INSTRUMENT

The GINA neutron reflectometer is a constant-energy angle-dispersive instrument with a horizontal scattering plane.¹⁸ The monochromator assembly, which is mounted in a gap of the cold neutron guide 10/3, selects the wavelength of the monochromatic beam within the range of 3.2 Å - 5.7 Å and can focus the beam in the vertical plane. In order to produce a polarized neutron beam,



FIG. 2. The layout of the GINA neutron reflectometer. The monochromator assembly (MONO) is mounted behind the concrete shielding (SH) on a turntable connected to the optical bench (B1) supporting the beam shutter (BS), (monitored by the semaphore control light (CL)), the intensity monitor detector (IMON), the cryostat of the beryllium filter (BF), the slit S1, the magnetic supermirror polarizer (P1), the adiabatic RF spin flipper (SF1) and the slit (S2). The bench is fixed to the central goniometer tower (ST), the components of which define the position and orientation of the sample surface relative to the beam and supports the components of the sample environment. The optical bench (B2) is connected to the turn-table underneath the central goniometer tower and it supports the slit (S3), spin flipper (SF2), the magnetic supermirror spin analyzer (P2), and slit (S4) in front of the detector unit (DET).

a magnetized polarizing supermirror (PSM) in transmission geometry and an adiabatic radio frequency (RF) spin flipper^{19,20} are used. The beam scattered by the sample undergoes spin analysis by an identical setup of a spin flipper and a spin filter, and finally it is detected by a two-dimensional position sensitive detector. To reduce the background, the detector is encased in a B₂O₃-mixed polyethylene shielding and four motorized Cd-blade slits are provided in the beam path to discard the undesired neutrons, including those scattered from the device components. The complete instrument setup is shown in Fig. 2. The neutron-optical devices of the reflectometer are mounted on two X95 optical benches²¹ to provide accurate definition of the beam height and a heavy-load support for various additional elements.

The optical bench B1 (cf. Fig. 2) defines the horizontal optical axis of the reflectometer and supports the beam shutter, the intensity monitor detector, the cryostat of the beryllium filter BF, the slit S1, the polarizer P1, the spin flipper SF1 and the slit S2. Bypassing the radiation shielding by a U-shape construction, the bench B1 is connected to the turntable of the monochromator, a heavy-load goniometer with vertical axis. The axis of the turntable coincides with that of the monochromator. The angle of the optical axis relative to the guide and consequently the wavelength can be changed by manually rotating the entire GINA setup around the turntable while air pads are activated and the entire setup floats over the marble floor. The allowed wavelengths are restricted at present to 3.2, 3.9, 4.6, 5.2 and 5.7 Å by the respective channels through the cylindrical concrete shielding



FIG. 3. Neutron spectra of regions of the monochromator crystals viewed by the pinhole-camera technique through the GINA monochromator (a), and aside from it (b). Dips at 4.6 Å and 2.3 Å are the basic and the first harmonic deflections by the monochromator graphite crystals. The difference between curves (a) and (b) are due to the neutron guide's Ni/Ti supermirror coating being viewed at different angles through the pinhole. Lower inset shows the detector image summed for the $4 \div 5$ Å interval. The upper inset is a horizontal divergence vs. wavelength map summed for the height of the summation areas (squares with sides of 29 mm) in the lower inset.

around the monochromator unit.

The downstream end of the optical bench B1 is fixed to the central goniometer tower ST, the components of which define the position and orientation of the sample surface relative to the beam and supports the various sample environment components (electromagnet, cryostat, etc.). The X95 optical bench B2 – the 2θ -arm of the reflectometer – is connected to the sample turntable underneath the central goniometer tower and it supports the slit S3, the spin flipper SF2, the spin analyzer P2, and the slit S4 in front of the detector unit along with its electronics and dedicated control PC. The encodercontrolled precise motion around the turntable is performed by a rubber-coated motorized wheel under the respective while the air pads are pressurized (cf. Fig. 2).

III. WAVELENGTH SELECTION AND FOCUSING

The monochromator is located five meters downstream from the cold source. The curved Ni/Ti supermirror neutron guide between the cold source and the GINA monochromator has a horizontal radius of 340 m, a height of 100 mm and a width of 25 mm. The inner, outer, top and bottom mirrors are of m = 2, 3, 2 and 2, respectively. The TOF neutron spectrum of the beam



FIG. 4. The monochromator assembly of the GINA reflectometer. Besides rotation, tilt and translation of the central graphite crystal (i.e. the full monochromator), the top and bottom pairs of graphite crystals can be individually rotated and tilted relative to it.

leaving this guide section was measured by the pinhole camera technique²² and the results are shown in Fig. 3. The parameters of the TOF experiment - monochromator to pinhole distance 3100 mm, pinhole to detector distance (flight length) 3750 mm, pinhole diameter: 3 mm, chopper open time: 0.1 ms, bin time: 8 μ s, detector gas absorption depth: 30 mm – resulted in a wavelength resolution²³ of $\Delta\lambda/\lambda = 0.01 + 0.12$ Å/ λ , leading to $\Delta\lambda/\lambda = 3.6\%$ at 4.6 Å and 6.2% at 2.3 Å, about 3 to 4 times worse than the actual monochromator resolution. Consequently, the dips appear wider and shallower in Fig. 3 at 4.6 Å and 2.3 Å than they are in reality. The monochromator assembly is mounted on a turntable connected to the optical bench B1 and comprises five highly oriented pyrolytic graphite (HOPG) crystals of $20 \times 20 \times 2$ mm³, attached to thin horizontal Al alloy rods (Fig. 4) and are lined up around the vertical axis.

The mounting rod of the central graphite crystal is directly attached to the monochromator bench, which is alined with respect to the axis of the monochromator turntable and can be rotated around its vertical axis. The rods of the remaining four crystals (two above and two below) are attached to small motorized 2-axis goniometers for horizontal alignment and vertical focusing. Initial alignment and focusing of the graphite crystals was facilitated by using a small chopper with a 1 mm diameter pinhole and a 2D detector in TOF mode to visualize the crystals (Fig. 5). The wavelength calibration was performed by fitting a harmonic series to the peak positions in the TOF spectrum of the central crystal (see the spectrum without the Be filter in Fig. 6a) thus calibrating the time channels in wavelength units. Vertical focusing to the sample position resulted in doubling the



FIG. 5. Intensity distribution (in arbitrary units) of the reflections of the monochromator crystals focused onto the sample position (2.5 m from the crystals) viewed at a distance of 2 m from a diaphragm of 1 mm diameter postioned in the focal point and recorded by the time-of-flight technique in different wavelength ranges of a width of 0.03 Å starting from 4.46 Å - 4.49 Å (a) to 4.68 Å - 4.71 Å (i), respectively. The differences in brightness of the crystals is either due to their relative misorientations and/or the wavelength-and divergence-dependent intensity distribution in the neutron guide.

intensity reflected by a $2\times2~{\rm cm^2}$ sample at grazing incidence as compared to the non-focused case of parallel graphite crystals.

The entire monochromator assembly can be tilted around a horizontal axis. When the instrument is not in operation, the monochromator assembly can be moved out of the beam using a linear x-stage. All eleven motions mentioned above are motorized and remotely controllable.

For the purpose of normalizing the measured reflectivity to the incident intensity, a low efficiency ${}^{3}\text{He}/\text{CF}_{4}$ beam monitor detector²⁴ with an active area of $H \times W =$ $100 \times 42 \text{ mm}^{2}$ is mounted in the beam path downstream of the monochromator assembly. The detecting efficiency of the intensity monitor is $\approx 0.1\%$ for $\lambda = 4.6$ Å neutrons.

Bragg reflections by the HOPG crystals contain higher harmonics according to the energy distribution of the incident neutron beam. These fractional-wavelength neutrons need to be filtered out from the beam incident onto the sample. GINA is equipped with a Be filter with Beslab size of (height×width×depth) 72 × 42 × 150 mm³ which is cold-finger-cooled by liquid nitrogen during regular operation. The transmission of the filter was measured using the TOF technique with a wavelength resolution of 6%. This experiment revealed that the filter has a transmission of 40.7% and 86.6% for $\lambda = 4.6$ Å, without and with liquid nitrogen cooling, respectively, while suppressing all higher harmonics at both temperatures (Fig. 6).

IV. SPIN POLARIZATION AND ANALYSIS

In order to produce polarized neutrons an Fe-Co/Si polarizing magnetic supermirror is used. The vertically oriented supermirror is mounted onto a rotator and a translator for adjustment to optimum polarization efficiency (defined below). The supermirror is placed between yokes of a permanent magnet construction with an in-plane vertical magnetic field of 30 mT. For spin analysis of the reflected beam, a PSM analyzer is used in an identical construction with the polarizer. By this setup, spin analysis of specular scattering is easily possible. In the case of off-specular scattering detailed studies are also possible but the experiments are rather time consuming. Both supermirrors are used in transmission geometry. The corresponding overall flipping efficiency (of P1, SF1 and P2) is shown in Fig. 7 as a function of



FIG. 6. A rea-integrated TOF spectrum of the detector pictures shown in Fig. 5 taken in the parallel-aligned orientations of the HOPG crystals without Be filter (a), and with Be filter at room temperature (b) and with liquid nitrogen cooling (c). The higher harmonics are suppressed by the filter at both temperatures. The transmission of the Be-filter is 40.7% and 86.6% for $\lambda = 4.6$ Å at room temperature and in the cold state, respectively.



FIG. 7. Polarization efficiency of the GINA setup vs. the incidence angle on the PSM of the analyzer P2, with SF1 ON and OFF, respectively, while SF2 was kept OFF if both cases. The overall polarization efficiency is ~ 0.9 in the optimum angular range of operation which is marked by the rectangle.

the position along the beam path.

In neutron reflectometry the signal to background ratio is always a critical parameter and has to be maximized. Suppression of scattering of neutrons by the beam-line components is the key issue. Therefore, instead of using Mezei flippers of simpler construction (always wires in the beam) we decided to opt for adiabatic RF spin flippers.^{19,20} The flipper coil is placed in a transverse static magnetic field with longitudinal gradient, produced by two iron plates energized by Nd-Fe-B permanent magnet stacks upstream and shunted by soft iron rods downstream. The RF coil for longitudinal RF field is part of a serial electric resonant circuit, with a sinusoidal current and bandwidth (full width at half maximum, FWHM) of $I_{\rm eff}$ = 4 A and 4.5 kHz at the resonance frequency of 175 kHz. The RF current is provided by a remote controlled power supply.²⁶ The parameters of the adiabatic RF spin flippers are summarized in Table 1. The flipper efficiency is better than 99% at any wavelength above 2 Å. The present flipper design is insensitive to the external field variations and/or influence of magnetic components along the beam path. It can be used in the entire wavelength range of GINA without further adjustment. The static field and the simulated $projection^{25}$ of the neutrons' magnetic moment onto the direction of the field are shown in Fig. 8 as a function of the neutron position along the beam path.

For neutron spin analysis an identical set of adiabatic RF spin flipper and PSM is placed downstream to the sample. The overall polarization efficiency of the setup was determined by measuring the reflected and transmitted intensity on the analyzer P2 with and without activating the spin flipper SF1, while the second flipper was kept OFF.



FIG. 8. Simulation of operation of the adiabatic radiofrequency spin flipper of GINA, displaying the measured static vertical magnetic field, $B_z(x)$, (left scale, curve (a)) and the simulated evolution of the z-component of the neutron spin $\langle S_z \rangle$ (right scale) along the coil axis, x for spin-up and spindown neutrons with wavelengths of 5.5 Å (b) and 3.2 Å (c), respectively. Vertical short-dashed lines indicate the coil ends. The parameters of the simulation are summarized in Table I.

The polarization efficiency was calculated by the formula

$$P = \frac{I^+ - I^-}{I^+ + I^-} \tag{1}$$

where I^+ and I^- are the corresponding reflected and transmitted intensities by the analyzer P2. This yields an overall efficiency of 0.9 for P1, SF1 and P2.

TABLE I. Parameters of the adiabatic RF spin flippers

Parameter	Range
beam height	60 mm
wavelength range	$(3.1 \div 5.8)$ Å
frequency	$175 \mathrm{~kHz}$
coil length / coil diameter	$300~\mathrm{mm}$ / $60~\mathrm{mm}$
number of turns	100
inductance / capacitance	$95~\mu\mathrm{H}$ / $8.8~\mathrm{nF}$
voltage / current (effective)	40 V / 4 A
magnet block (Nd-Fe-B) $2\times$	$60\times50\times20~\mathrm{mm^3}$
yoke $(L \times W \times H) \ 2 \times$	$500\times100\times12~\mathrm{mm}^3$
magnetic field in the centre	$5.6 \mathrm{mT}$
longitudinal gradient	$(0.2 \div 0.4) \text{ mT/cm}$

V. SAMPLE POSITIONING

The flat sample is mounted on an adjustable vertical flat surface attached to the top cradle of the central goniometer tower. The sample mounting depends on the sample environment. For room temperature reflectivity measurements the flat surface has a small bore through which the sample is sucked to the vertical surface and held in position during the experiment by a small vacuum pump. In such a way undesired scattering by the fixing elements is minimized. Symmetrical sample positioning is ensured by using two cradles and two perpendicular linear stages. The cradles and translators position the sample in the vertical plane and set the sample surface orientation. The θ and 2θ angles are encoder controlled for increased precision. Fine positioning of the beam is maintained by several slits with cadmium blades. The slits can be opened in the range of 0 to 10 mm with a precision better than 0.2 mm. One slit (S1) is placed downstream of the Be filter and one (S2) downstream of flipper SF1, just upstream of the sample. The slit S1 defines the beam on the polarizer mirror to decrease the beam divergence thus to increase the polarization ratio. Slit S2 decreases the beam divergence on the sample and absorbs the neutrons which might not reach the sample or scattered off by the polarizer. With these optical elements the setup can achieve a relative Q-resolution of 10% to 2% for the available Q-range of 0.005 to 0.25 Å⁻¹.

VI. SAMPLE ENVIRONMENT

GINA is primarily dedicated to reflectometry of magnetic heterostructures. For studies of magnetism, vital environmental parameters are (low) temperature and (occasionally high) applied magnetic field. A closed cycle ⁴He cryostat (comprising a cold finger setup²⁷ and evacuated by a small turbomolecular pump) can be mounted on the central goniometer tower of GINA with or without the electromagnet. The sample temperature can be varied in the nominal 9 to 300 K range.²⁸ The GINA beam line is equipped with an air-cooled electromagnet which produces magnetic fields up to 0.55 T for the pole distance of 40 mm that accommodates the 1.5" diameter Al cap of the cryostat. A water-cooled air core coil pair provides smaller magnetic fields up to approx. 35 mT.

VII. NEUTRON DETECTION AND BACKGROUND REDUCTION

For detecting the neutrons, a multi-wire proportional chamber filled with ${}^{3}\text{He}/\text{CF}_{4}$ gas mixture of 2.5/3 bar partial pressures with 200 × 200 mm² active area and spatial resolution of 1.6 mm (FWHM). In order to suppress the background, the detector is encased in a polyethylene shielding of 30 mm thickness containing 20 wt% natural B₂O₃. The two-dimensional spatial detection is managed by two delay lines and the positions are determined by a DASY TDC module (produced by ESRF, Grenoble) installed in a slot of the detector PC which is dedicated exclusively to the detector data acquisition and mounted on the 2 θ -arm of the reflectometer. When detecting specular scattering, two slits (S3 and S4) are placed in front of the detector window to discard undesired radiation. If no spin analysis is required, for further background suppression, an evacuated flight tube is mounted along the entire length of the 2θ detector arm. Mounting the spin analyzer and flipper in a vacuum vessel is a plan for the future.

VIII. INSTRUMENT CONTROL

The GINA hardware and the control software are rather flexible and are designed for maximum remote controllability. In its full configuration, GINA comprises more than 30 stepping motors.²⁹ Certain critical motions, such as θ - and 2θ - angles and precision slit positions are absolute or relative encoder-controlled. A custom made unit built around a USB multi-function data acquisition module³⁰ controls the air compressor, the air pads, the liquid nitrogen level and temperature in the Be-filter, the beam shutter and its control lights, the beam intensity monitor and the various modular DC power supplies (the latter ones to energize electromagnets, and various coils in the setup including optional Mezei flippers). The high voltage power supplies of the detector and that of the beam monitor, the linear amplifiers, the discriminators and the ratemeters are realized in modules of NIM standard. The control PC directly communicates with the detector PC via ethernet and with the indexer modules of the motion control units as well as with the temperature controller via RS232 lines. All listed components are controlled by the GINASoft control software written in LabView V10.0 for MS Windows. The user interface of the program allows for various alignment and scan modes as well as changing polarization and sample environment (flipper current and frequency, temperature, magnet current, etc.) remotely. For increased user comfort the command format of the user interface is user-configurable, and includes a command format that resembles to that of SPEC.³¹ 2D detector pictures and reduced reflectivity data can be efficiently viewed and manipulated during data acquisition. Collected and manipulated data as well as extended log information (including graphics) are saved in a clearly structured database format. Human control is facilitated by a web camera which is installed in the control PC. Using remote desktop option, most operations can be performed remotely via internet from outside the experimental hall or even from a distant continent.

IX. DATA HANDLING AND EVALUATION

Users of the GINA reflectometer are offered the data handling and evaluation software FitSuite,³² a thoroughly documented program with a detailed project home page, written for Windows *and* Linux, which is presently suitable for evaluating data of 14 experimen-

TABLE II. Operation parameters of the GINA neutron reflectometer at the Budapest Neutron Centre

Parameter	Range
wavelength	$(3.2 \div 5.7)$ Å in five steps
present wavelength	4.6 Å
max. scattering angle (θ)	$\geq 35^{\circ}$
angular resolution $(\Delta \theta)$	0.003°
$\Delta\lambda/\lambda$	$\sim 1\%$
background level at the 2D de-	0.01 cps cm^{-2}
tector	
detector	$\rm 2D~PSD~200\times 200~mm^2$
detector resolution	$1.6\times1.6~\mathrm{mm^2}$
neutron flux at the monochro-	$4\times 10^5~\rm{n}\times \rm{cm}^{-2}\times \rm{s}^{-1}$
mator position	
background reflectivity	$< 10^{-5}$
overall polarization efficiency of	0.9
polarizer and analyser SM	

tal methods, including specular polarized neutron reflectometry (with e.g. model for diffusion, also in isotopeperiodic multilayers) and off-specular (diffuse) polarized neutron reflectometry (in the distorted-wave Born approximation) as well as specular x-ray reflectometry. Specular and off-specular reflectivities are calculated using the matrix approach Ref. 3 and the Distorted Wave Born Approximation given in Ref. 33. FitSuite³² handles the corresponding theories and sample structures consistently in a common structure and allows for parameter restrictions, correlations and simultaneous simulation and fit of models to the experimental data.

X. EXAMPLE REFLECTOGRAMS

Two example reflectograms are chosen to highlight the present performance of the GINA setup. The first one is the specular reflectivity of a 4-inch Si wafer measured in non-polarized mode. The data shown in Fig. 9 were collected over a time of 22 hours. The fit to the experimental points gave an SLD value of $(3.03 \pm 0.03) \times 10^{-6} \text{ Å}^{-2}$ for the native surface oxid layer of 20 Å thickness and roughnesses of 4 A for both the silicon and the native oxide surfaces, respectively. The SLD for the silicon was kept constant to the nominal value of 2.07×10^{-6} Å⁻². The obtained SLD value for the oxide layer is slightly smaller than the SLD value of 3.47×10^{-6} Å⁻² for amorphous bulk SiO_2 . This deviation may result from the lower density of the thin oxide layer as compared to that of the bulk. The fit, full line in Fig. 9, provides a background reflectivity of 9×10^{-6} .

The second example is a magnetized Ni film. Polarized neutron reflectometry provides a means to simultaneously determine the atomic layer and magnetization profile in a multilayer. Deviation of both atomic and magnetic momentum density of Ni films from the bulk values



FIG. 9. The specular reflectivity curve of a four-inch Si wafer measured at the GINA reflectometer in non-polarized mode. The fitted background reflectivity is 9×10^{-6} .

was a subject of earlier neutron reflectometric studies. Singh and Basu studied a Ni film of 1500 Å on glass substrate (a neutron mirror) by unpolarized and polarized neutron reflectometry.^{34,35} Their analysis of the data provided a nearly 50% decrease of the atomic and a similar extent of decrease in the magnetic momentum density in a surface layer of 235 Å and an about 10% decrease of both quantities in the rest of the film without any chemical change. Therefore we decided to study such effects on a carefully prepared Ni film. In order to exclude undesired surface effects, an isotopic multilayer Ni film of nominal layer structure $MgO/[^{62}Ni(5nm)/^{nat}Ni(15nm)]_5$ was prepared by molecular beam epitaxy in the Wigner Research Centre for Physics, Budapest depositing Ni of natural isotopic abundance by electron beam evaporation and 96.2% enriched 62 Ni from a Knudsen cell onto (100) MgO of $20 \times 20 \times 2 \text{ mm}^3$ size. Reflection high energy electron deposition images taken during deposition revealed both epitaxial and polycrystalline regions in the multilayer. Due to the large negative SLD of 62 Ni, it provides a sensitive probe of any structural change in the multilayer. Specular reflectivity of the sample was measured on the GINA reflectometer at room temperature in polarized mode without polarization analysis. Prior to the experiment the sample was magnetized from the virgin state by 50 mT in-plane external magnetic field parallel to the guide field. The total data collection time was 56 hours. In Figs. 10a–10b the reflectivities R^+ and R^- are shown together with the curves of the simultaneous fit by $FitSuite^{32}$ (latter in full lines). For clarity, the measured and fitted spin asymmetry $(R^+ - R^-) / (R^+ + R^-)$ curves are displayed separately in Fig. 10c. The simultaneous fit was constrained to a periodic layer structure and yielded layer thicknesses of (175 ± 5) Å and (53.5 ± 5) Å for ^{nat}Ni and ⁶²Ni, respectively and common rms interface roughness of (5 ± 2) Å. The errors are 1σ statistical errors as obtained



FIG. 10. Measured spin-up R^+ , and spin-down R^- reflectivities and the calculated $(R^+ - R^-)/(R^+ + R^-)$ spin asymmetry of the isotope-periodic multilayer MgO (001) /[⁶²Ni(5 nm)/^{nat}Ni(15 nm)]₅ measured on the GINA reflectometer in polarized mode without polarization analysis.

by the least-squares fit. The fitted scattering length densities were: $\text{SLD}(^{\text{nat}}\text{Ni}) = (9.13 \pm 0.5) \times 10^{-6} \text{ Å}^{-2};$ $SLD(^{62}Ni) = (-7.0 \pm 1) \times 10^{-6} \text{ Å}^{-2}$ showing only minor deviations from their known bulk SLD values, consequently as well as from the bulk Ni atomic density. The magnetization in the ^{nat}Ni and ⁶²Ni layers were kept identical in the fit which provided (0.44 ± 0.12) T. This value amounts only to 66% of the known room temperature saturation magnetization MS of the bulk Ni $(55 \text{ emu/g})^{36}$ corresponding to 0.67 T). This may be either due to a decrease of the Ni magnetic moment in the layers or an incomplete magnetic saturation of the sample. The reflectivity measurement was performed in 50 mT external magnetic field. This in-plane field strength was considered sufficient for saturation, since the bulk Ni saturates at 300 Oe even along the (100)direction of the hard magnetization.³⁶

However, the magnetooptical Kerr-loop (Fig. 11) taken subsequent to the reflectivity measurement revealed a partial (74%) saturation of the Ni film in 50 mT. The somewhat lower value of 66% provided by the reflectivity fit can be explained by a lower magnetization of the near-edge regions of the sample, contributing to the re-



FIG. 11. Magnetooptical Kerr loop of the Ni isotope multilayer sample showing an incomplete magnetic saturation of the layer at 50 mT, the external field used in the reflectivity measurement. Inset: Scanning electron microscope picture of a 200 × 130 μ m surface region of the sample. Tubular discontinuities amount up to about 1% of the film volume.

flectivity but not to the MOKE signal. Scanning electron microscopic pictures of the sample surface (inset in Fig. 11) reveals sparse plane-perpendicular tubular discontinuities of the multilayer. Image analysis provides 1.37% surface coverage of the tube openings. Therefore the total volume of the tubes and their effect on the average film density is negligible in the present sample.

XI. SUMMARY

We have shown that GINA, the newly installed dancefloor-type constant energy angle-dispersive neutron reflectometer at the Budapest Neutron Centre is a versatile instrument in both polarized and unpolarized modes of operation. Examples were given for both modes of operation. The sample orientation of the reflectometer is vertical. The available sample environment facilities are: closed cycle cryostat optionally combined with external magnetic field by an iron core electromagnet up to 0.55 T or air core coils for 35 mT in various directions, spin polarization and polarization analysis by single polarizing supermirrors. Detection of specular and diffuse scattering is facilitated by a two-dimensional positionsensitive detector. All components of the instrument are controlled by a program written in LABVIEW. The program allows for alignment and scan modes as well as changing polarization and sample environment parameters remotely. Reflectivities above five orders of magnitude have been measured with further improvements underway. Further developments including an environmental cell for biomimetic membrane studies, an electromagnet with higher fields and orientation versatility and a supermirror fan analyzer and further background

suppression elements are also planned. The GINA reflectometer is open for Hungarian and international users throughout the year.³⁷

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